Synchronization of coupled oscillators

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The Kuramoto model

The equation

$$\dot{\theta_j} = \omega_j + \frac{\kappa}{N} \sum_{k=1}^N \sin(\theta_k - \theta_j), \quad j = 1, \dots, N$$

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is the Kuramoto model with parameters

- ω_j : frequency of each oscillator
- K: coupling strength
- N: number of oscillators

 θ_j is a phase variable in $[0, 2\pi]$.

Initial conditions

Choose N random frequencies from a distribution g(ω)
Typically, g(ω) is unimodal (Gaussian or Lorentzian)
Choose initial phases (e.g. randomly, or constant, θ_i = 0, ∀j)

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All-to-all coupling: $\mathcal{O}(N^2)$ Local coupling on a circle "Small world networks" — random couplings with given probability

Further variations:

Delayed couplings (neural nets)

Different populations of oscillators (different natural frequencies) External driving: periodic, quasi-periodic, noise, etc.

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Increasing K

As the coupling increases, the oscillators synchronize.



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The order parameter

Phases are taken modulo 2π . Use *circular statistics*.

$$re^{i\psi}=rac{1}{N}\sum_{k=1}^{N}e^{i heta_{k}}$$

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Order parameter: a complex number in polar form

 $\psi \in [0, 2\pi]$ is the "average phase" $r \in [0, 1]$ is the phase coherence Oscillators in sync $\rightarrow r \approx 1$ independent / unordered $\rightarrow r \approx 0$

Order parameter geometry

The order parameter is a vector. Its length increases as oscillators synchronize.



Red arrow: order parameter Blue circles: phase values

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Onset of sync — critical coupling

For $K < K_c, r \approx 0$ For $K > K_c, r$ approaches 1.



 K_c depends on the spread of the frequency distribution $g(\omega)$

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Time evolution of the order parameter

Coupling dependece: For $K > K_c$ the order parameter rapidly increases.



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The Kuramoto model is studied in the limit $N \to \infty$. Continuous probability distribution $\rho(\theta)$ of phase A partial differential equation describes the evolution of $\rho(\theta)$ Variable substitutions: The order parameter is explicitly used Rescaling of variables: let $\langle g(\omega) \rangle = 0$ The **Ott-Antonsen ansatz** (2008): the dynamics of the Kuramoto model can be reduced to a low-dimensional ODE.

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Phenomena in coupled oscillators

Phase locking

- Incoherent to synchronized transition
- Oscillation death
- Chimera states
- Noise-induced synchronisation

Examples:

phase locking and oscillation death in organ pipes period doubling in applause (Néda et al. 2000)

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Landau-Stewart oscillators

The Hopf oscillator (or Landau-Stewart oscillator):

$$\dot{z} = (\alpha + i\omega - |z|^2)z$$

for complex z.

Limit cycle oscillator with amplitude $\sqrt{\alpha}$ if $\alpha > 0$, and frequency ω .

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Coupled Landau-Stewart oscillators

Simplest case of mean field coupling:

$$\dot{z}_j = (\alpha + i\omega_j - |z_j|^2)z_j + \frac{\kappa}{N}\sum_{k=1}^N z_k, \quad j = 1, \dots, N$$

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K < 0 is inhibitory K > 0 increases amplitude Reduces to the Kuramoto model if the amplitude is neglected (Pyragas et al. 2007).

Order parameter

Define the order parameter for the Landau-Stewart oscillator

$$re^{i\psi} = rac{1}{N}\sum_{k=1}^{N}e^{i\arg z_k}$$



Adaptive oscillators for beat tracking uses external forcing, e.g.

$$\dot{z}_j = f(z_j) + K\bar{Z} + x(t)$$

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f(z) is the Landau-Stewart oscillator \overline{Z} is the average of oscillators x(t) is the input signal

Sync in chaotic systems

Discovered by Pecora & Carroll (1990)

$$\begin{aligned} \dot{\mathbf{x}} &= \mathbf{f}(\mathbf{x}) + \mathcal{K} y_1 \\ \dot{\mathbf{y}} &= \mathbf{f}(\mathbf{y}) + \mathcal{K} x_1 \end{aligned}$$

where $\mathbf{x}, \mathbf{y} \in \mathbb{R}^3$ or higher dimension. Criterion of sync: $\|\mathbf{x} - \mathbf{y}\| < \epsilon$ Generalized synchronization (of phase): $\dot{\theta}_1 = \dot{\theta}_2$ — same frequency, phase may differ.

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Further reading

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